

APPENDIX

1

DESIGN NOTE

A compact low-voltage TEA-N₂ laser

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Abstract. A compact TEA-N₂ laser has been designed using a double-fold, two-stage Blumlein circuit which enabled operation at low voltage (4.5–7.5 kV) and a laser channel gap of 1 mm. An output energy of 150–370 μ J per pulse was obtained and its average electrical-to-optical conversion efficiency of about 0.06% was comparable to the values reported for most compact TEA-N₂ lasers operated at 15–40 kV.

Keywords: TEA-N₂ laser, double-fold, two-stage Blumlein, low-voltage laser

1. Introduction

A compact, transversely excited atmospheric-pressure nitrogen (TEA-N₂) laser with a few hundred microjoules of output energy provides a convenient ultraviolet (UV) light source for dye-laser pumping, fluorescence lifetime studies, and optical diagnostics of transient plasmas. Although several compact designs of parallel-plate, Blumlein-driven TEA-N₂ laser have been reported [1–5], their operating voltages are usually in the range of 15–40 kV. It is always desirable to lower the operating voltage so as to reduce the generation of electromagnetic interference and to minimize the dielectric breakdown of the insulator layer in parallel-plate capacitors. A TEA-N₂ laser was reported [5] to operate at 5–10 kV, but its output energy of 40–150 μ J was barely sufficient for most of the above applications. A miniature TEA-N₂ laser employing 3 cm silicon electrodes was recently reported [6] to operate at 5 kV, but its laser output energy was merely 3 μ J per pulse.

In previous articles [7, 8], we reported on the improved performance of a TEA-N₂ laser using a two-stage Blumlein circuit. Although the input power is significantly improved in a two-stage Blumlein circuit owing to an increase of the breakdown voltage for laser discharge initiation by about 25%, the total size of the parallel-plate capacitors is increased several times [1–5]. We thus attempted to reduce the total capacitor size by using both double-fold and overlay techniques for packaging the TEA-N₂ laser system. To enable low-voltage operation, the UV preionizers were geometrically optimized to prevent arc formation in the laser discharge and a swinging cascade spark gap of lower inductance was used. As a result, a compact TEA-N₂ laser, of dimensions 25 \times 17 cm², was operated at 4.5–7.5 kV and delivered an output energy of 150–370 μ J per pulse.

2. Experimental set-up

A laser channel of length 25 cm and gap 1 mm was formed between two laser electrodes, each of which had a cross section of 4 \times 4 mm². These electrodes, as shown in figure 1(a), were shaped out of the edge of two brass plates 25 cm long, 6 cm wide and 1.2 cm thick. A gas flow tunnel, with a cross section of 3 \times 3 mm², was milled at the edge of the two brass plates. N₂ gas was flushed into the laser channel via equally spaced 1 mm² transverse grooves along the two flow tunnels. The laser channel was sealed by an aluminium-coated back mirror at one end and a microscope slide at the output end. The back mirror, however, could only increase the laser output by about 20% owing to the narrow laser channel and lack of a mirror alignment mechanism. Two UV preionizers were introduced above and below the laser channel; UV radiation was provided by corona discharges between the two opposite aluminium strips protruding beneath the laser electrodes. In our earlier work [7], the position of these aluminium strips was not critical to the laser output consistency since the operating voltage was \geq 10 kV. In this work, where the operating voltage was between 4.5–7.5 kV, arc discharges occurred readily in the laser channel at the initial development stage. These were strongly suppressed by aligning the edges of two aluminium foils with the laser channel, one of which was directly above and the other one directly below (see figure 1(a)). An optimization study of the position of these aluminium foils could not be carried out since these foils were totally obscured by the various layers of Mylar, aluminium foil and foam.

A double-fold, two-stage Blumlein circuit (figure 1(b)) was used, which consisted of parallel-plate capacitors on both sides of the brass plates. Thus, the two capacitors responsible for the laser discharge, C₁ and C₂, were formed between the inner aluminium foil, Al-1, and the brass plates, L and RP respectively. The dielectric insulator used for C

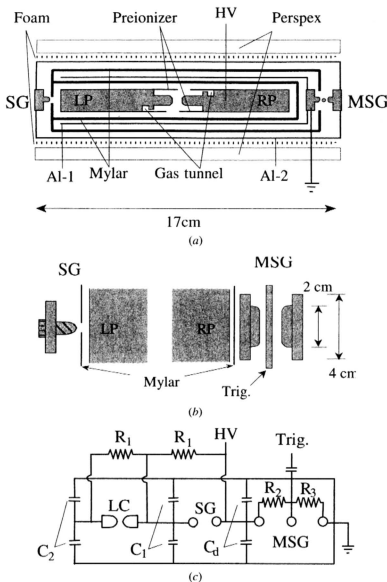


Figure 1. (a) End-on view of the double-fold, two-stage Blumlein-driven TEA-N₂ laser: MSG, main spark gap; SG, second spark gap. (b) Top view of the MSG and the SG. (c) Equivalent circuit of the TEA-N₂ laser. R_1 , R_2 , R_3 are 100 M Ω , 40 M Ω and 60 M Ω respectively. $C_1 = C_2 = 3.2$ nF and $C_d = 6.4$ nF; LC is the laser channel with gap varying between 1–1.5 mm.

and C_2 was formed from two layers of 50 μ m Mylar. By overlaying Al-1 with two layers of 50 μ m Mylar followed by another aluminium foil, Al-2, the dummy capacitor, C_d , was formed. The laser assembly was packaged by pressing the multilayer structures between two thin layers of foam and two 1.2 cm thick Perspex plates. These capacitances were estimated, using $C = \epsilon A/d$, to be $C_1 = C_2 = 3 \times 2$ nF and $C_d = 6 \times 2$ nF.

The main spark gap (MSG), shown in figure 1(b), consisted of three electrodes arranged in parallel. The centre electrode is a 1 mm rod used as a trigger pin, and the other two breakdown electrodes are semi-cylindrical rods. The MSG was tightened by two screws to the top and bottom Perspex plates. The charging voltage across the MSG was resistively divided in a ratio of $R_2:R_3 = 1:2$ between the two breakdown gaps: the first between the high-voltage anode and the centre electrode (1 mm pin) and the second between the centre and ground electrodes. Upon triggering at the centre electrode, the first gap would break down first which

then triggered a breakdown across the second gap to the ground electrode. The main reason for using this flat-type of swinging cascade spark gap instead of a normal trigger was due to the difficulty in introducing a triggering pin to the side of MSG chamber which provided easy access for the trigger pulse. In this design, the centre electrode was exposed to the side of MSG chamber which provided easy access for the trigger pulse. The trigger pulse was obtained from a SC circuit the output pulse of which was 300 V and was stepped up to more than 12 kV by a car ignition coil.

3. Results and discussion

The laser output energy as a function of both the charging voltage, V_o , and the breakdown gap in the second spark gap (SG) is presented in figure 2. A minimum breakdown gap in the SG of about 0.3–0.5 mm was necessary to achieve a high laser output at any operating voltage from 4.5–7.5 kV. For a SG breakdown gap smaller than 0.3 mm

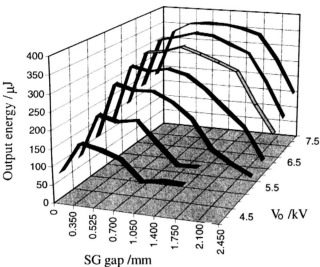


Figure 2. Laser output energy as a function of charging voltage (V_0) for different SG breakdown gaps.

the laser output energy not only decreased sharply, it was also irreproducible. This was caused mainly by arc formation in the laser discharge. As the SG gap was increased to more than 0.5 mm, the laser output energy began to decrease gradually and its shot-to-shot variation increased. These were caused mainly by the presence of streamers in the laser discharge which gave rise to a spatially nonuniform input power loading and optical inhomogeneity. The effects of the SG gap on the laser output curve were investigated previously [7].

The temporal sequences of spark gap breakdown in the MSG and the SG with respect to laser discharge are shown in figure 3, represented by their respective photodiode (FND 100, EG&G) signals of light emissions from the MSG and SG and the laser emission. These signals, which were measured with a SG gap of 0.5 mm and a charging voltage V_0 of 6 kV, were recorded using a 200 MHz digital oscilloscope (Teknatrix TDS 320). The time delay between MSG and SG firings was between 28–18 ns, which decreased linearly as V_0 was increased from 4.5–7.5 kV. The change in the time delay between the laser pulse and SG firing could not be measured accurately, although it is expected to decrease accordingly. It should be noted that the laser pulse was about 3 ns at its FWHM—this may not be the true laser pulse width but the risetime of the 200 MHz oscilloscope.

The risetime of the photodiode signal for the SG was about 5 ns, as compared with about 15 ns for the MSG. The difference is mainly due to a C_d capacitance in the first discharge loop, C_d-L_1 -ground, as shown in figure 1(c). On the other hand, when the SG was fired the second discharge loop was much faster because C_1 and C_d formed a reduced capacitance, $C_1 = C_1 C_d / (C_1 + C_d) = 4$ nF. The inductance, L_1 , of the MSG may be estimated by assuming that the timing of the laser pulse, as previously observed to correspond to the maximum voltage inversion in the first discharge loop [7], was at $3/4$ of the period, T , of the first discharge loop. Thus, $T = 2\pi(L_1 C_d) \approx 47$ ns and $L_1 \approx 4.7$ nH. The inductance, L_2 , of the SG may also be estimated as the laser pulse emerged around the maximum voltage inversion of the second discharge loop [7], thus $T = 2\pi(L_2 C_s) \approx 12 \times 2$ ns. Since $C_s = 4$ nF, $L_2 \approx 3.6$ nH.

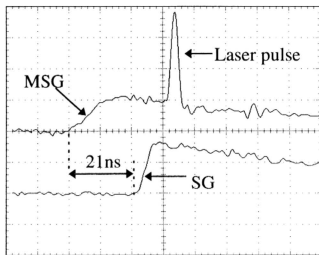


Figure 3. Temporal characteristics and time delays between light emissions from the MSG and the SG and the laser pulse.

The electrical-to-optical conversion efficiency, calculated by dividing the laser output energy by the total energy stored in the three capacitors ($1/2(C_1 + C_2 + C_d)V_0^2$), was on average close to 0.06%. Although this value is comparable to the reported values of around 0.05% of other compact systems operating at 15–40 kV [1, 4–6], it is lower than 0.18% reported in [3]. The high efficiency in [3] was mainly attributed to the back mirror which increased the laser output 2.5 times, as compared to an increase of only 20% in this work. If we neglect the feedback effect by the back mirror, the single-pass conversion efficiency of 0.05% obtained in this work is more comparable to the 0.072% reported in [3].

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4. Conclusion

A compact TEA- N_2 laser has been designed and operated at voltages from 4.5–7.5 kV to deliver an output energy of 150–370 μ J. Despite a large dummy capacitance, $C_d = 12$ nF, the electrical-to-optical conversion efficiency of about 0.06% is comparable to most of the reported values obtained at 15–40 kV. The laser was also tested at 4 kV, but its output energy in this case was less than 70 μ J and its shot-to-shot output variation was more than 25%.

One of the prerequisites of the present system is a narrow laser channel gap in order to enable low-voltage operation; this was reduced from the normal value of 3 mm to 1 mm. Another is alignment of the edges of two aluminium foils with the laser channel, one directly above and the other directly below the laser channel. This has hardly ever been reported to be important since the operating voltage of TEA- N_2 laser is seldom below 10 kV. However, in the recently reported miniature TEA- N_2 laser, it was noted that streamer or plasma filaments occurred readily at a 5 kV operating voltage [5]. Thus, our results may suggest that sufficient UV radiation at several kilovolts can still be harnessed from corona discharges along the edges of the aluminium foils.

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